

Numerical Modeling of Methane Hydrate Evolution

Nathan L. Gibson

Joint work with F. P. Medina, M. Peszynska, R. E. Showalter

Department of Mathematics



SIAM Annual Meeting 2013
Friday, July 12

*This work was partially supported by NSF DMS-1115827 "Hybrid Modeling in porous media".

Outline

1 Introduction

Outline

- 1 Introduction
- 2 Model Development
 - Conservation of mass
 - Solubility constraints

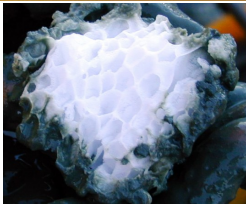
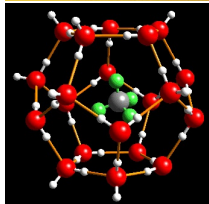
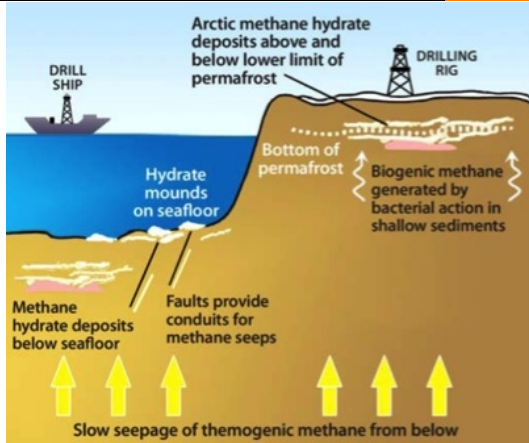
Outline

- 1 **Introduction**
- 2 **Model Development**
 - Conservation of mass
 - Solubility constraints
- 3 **Abstract Evolution Equation**
 - Examples
 - Analysis

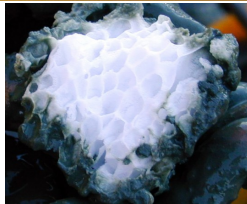
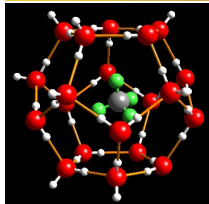
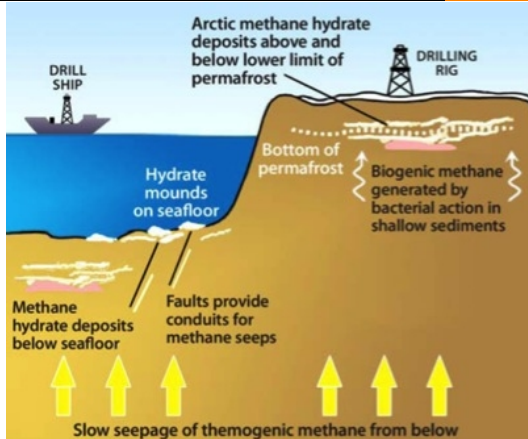
- 1 **Introduction**
- 2 **Model Development**
 - Conservation of mass
 - Solubility constraints
- 3 **Abstract Evolution Equation**
 - Examples
 - Analysis
- 4 **Numerical aspects**
 - Fully discrete scheme
 - Semismooth Newton solver

- 1 **Introduction**
- 2 **Model Development**
 - Conservation of mass
 - Solubility constraints
- 3 **Abstract Evolution Equation**
 - Examples
 - Analysis
- 4 **Numerical aspects**
 - Fully discrete scheme
 - Semismooth Newton solver
- 5 **Some experiments**

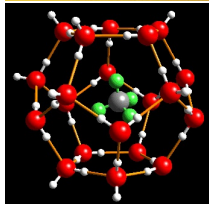
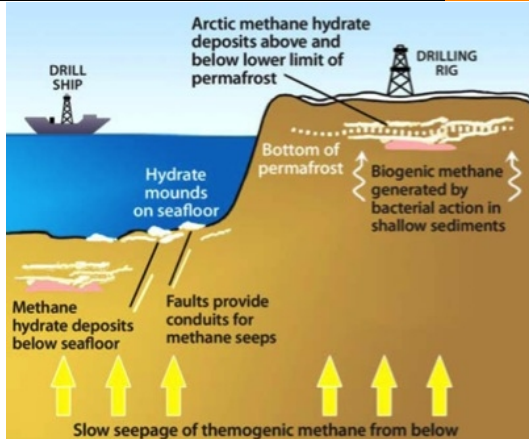
- 1 **Introduction**
- 2 **Model Development**
 - Conservation of mass
 - Solubility constraints
- 3 **Abstract Evolution Equation**
 - Examples
 - Analysis
- 4 **Numerical aspects**
 - Fully discrete scheme
 - Semismooth Newton solver
- 5 **Some experiments**
- 6 **Conclusions and future work**



- Methane hydrates are an ice-like substance containing methane molecules trapped in a lattice of water molecules.



- Methane hydrates are an ice-like substance containing methane molecules trapped in a lattice of water molecules.
- We will consider a simplified model of evolution of methane hydrates in the hydrate zone of the sea-bed.



- Methane hydrates are an ice-like substance containing methane molecules trapped in a lattice of water molecules.
- We will consider a simplified model of evolution of methane hydrates in the hydrate zone of the sea-bed.
- Components:
 - ▶ CH_4
 - ▶ H_2O

*[Images from DOE-NETL]

Conservation of mass for CH_4 component

Let $\Omega \subset \mathbb{R}^3$ be a bounded region of points $x \in \Omega$.

Parameters: (assumed given)

ϕ_0 porosity

K_0 permeability

f_M external source of CH_4

Conservation of mass for CH_4 component

Let $\Omega \subset \mathbb{R}^3$ be a bounded region of points $x \in \Omega$.

Parameters: (assumed given)

ϕ_0 porosity

K_0 permeability

f_M external source of CH_4

Assumptions:

- Pressure $P(x)$ given by hydrostatic gradients
- Temperature $T(x)$ given by geothermal gradients
- High pressure and low temperature imply Hydrate zone:
only liquid and hydrate phases present

Conservation of mass for CH_4 component (cont.)

Unknowns:

- S_l, S_h saturations (void fractions), with $S_h = 1 - S_l$
- $\chi_{IM}, \chi_{IW}, \chi_{IS}$, with $\chi_{IM} + \chi_{IW} + \chi_{IS} = 1$
mass fractions of methane, water and salt in liquid phase

Conservation of mass for CH_4 component (cont.)

Unknowns:

- S_l, S_h saturations (void fractions), with $S_h = 1 - S_l$
- $\chi_{IM}, \chi_{IW}, \chi_{IS}$, with $\chi_{IM} + \chi_{IW} + \chi_{IS} = 1$
mass fractions of methane, water and salt in liquid phase

Note:

- Salinity, χ_{IS} , assumed known and fixed to that of seawater
- χ_{hM}, χ_{hW} mass fractions in hydrate phase (assumed known)

Conservation of mass for CH_4 component (cont.)

Unknowns:

- S_l, S_h saturations (void fractions), with $S_h = 1 - S_l$
- $\chi_{IM}, \chi_{IW}, \chi_{IS}$, with $\chi_{IM} + \chi_{IW} + \chi_{IS} = 1$
mass fractions of methane, water and salt in liquid phase

Note:

- Salinity, χ_{IS} , assumed known and fixed to that of seawater
- χ_{hM}, χ_{hW} mass fractions in hydrate phase (assumed known)

Conservation of mass for CH_4 component

$$\frac{\partial}{\partial t} (\phi_0 S_l \rho_l \chi_{IM} + \phi_0 S_h \rho_h \chi_{hM}) - \nabla \cdot (D_{IM} \rho_l \nabla \chi_{IM}) = f_M \quad (1)$$

where densities ρ_l, ρ_h and diffusion D_{IM} are assumed constant.

Unified notation

We redefine parameters

$$R := \frac{\rho_h \chi_{hM}}{\rho_l}, f := \frac{f_M}{\rho_l \phi_0}, D_0 := \frac{D_{IM}}{\phi_0},$$

and redefine variables

$$S := S_l, v := \chi_{IM}, u := S v + R(1 - S),$$

so that (1) becomes

$$\frac{\partial u}{\partial t} - \nabla \cdot (D_0 \nabla v) = f, \quad (2)$$

where we may further scale the problem so that $D_0 = 1$.

Still need $\langle S_I, \chi_{IM} \rangle \in \mathcal{F}(x)$ provided by **solubility constraints**. First, we assume that the **maximum solubility constraint** $\chi_{IM}^{\max}(x)$ is given.

Still need $\langle S_i, \chi_{iM} \rangle \in \mathcal{F}(x)$ provided by **solubility constraints**. First, we assume that the **maximum solubility constraint** $\chi_{iM}^{\max}(x)$ is given.

Then the solubility satisfies a nonlinear complementarity constraint (NCC)

Solubility constraint

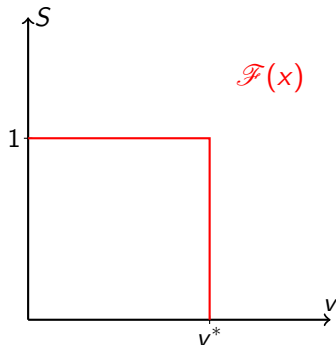
$$\begin{cases} \chi_{iM} \leq \chi_{iM}^{\max}, & S_i = 1, \\ \chi_{iM} = \chi_{iM}^{\max}, & S_i \leq 1, \\ (\chi_{iM}^{\max} - \chi_{iM})(1 - S_i) = 0. \end{cases} \quad (3)$$

Still need $\langle S_I, \chi_{IM} \rangle \in \mathcal{F}(x)$ provided by **solubility constraints**. First, we assume that the **maximum solubility constraint** $\chi_{IM}^{\max}(x)$ is given.

Then the solubility satisfies a nonlinear complementarity constraint (NCC)

Solubility constraint

$$\begin{cases} \chi_{IM} \leq \chi_{IM}^{\max}, & S_I = 1, \\ \chi_{IM} = \chi_{IM}^{\max}, & S_I \leq 1, \\ (\chi_{IM}^{\max} - \chi_{IM})(1 - S_I) = 0. \end{cases} \quad (3)$$



In our unified notation, we let $v^* = \chi_{IM}^{\max}$, so that

$$\langle v, S \rangle \in \mathcal{F}(x; \cdot) := [0, v^*(x)] \times \{1\} \cup \{v^*(x)\} \times (0, 1]$$

Or equivalently,

$$v \in \alpha_{MH}(x; u) := (u - v^*(x))_- + v^*(x), \quad u \leq R. \quad (4)$$

We consider the initial boundary-value problem

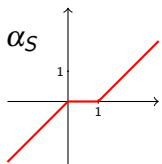
Abstract Evolution Equation

$$\begin{aligned} \frac{\partial u}{\partial t} - \Delta v &= f, & v &\in \alpha(u) \text{ on } \Omega \times (0, T) \\ v &= 0 & &\text{ on } \partial\Omega \times (0, T) \\ u(\cdot, 0) &= u_0(\cdot) & &\text{ on } \Omega. \end{aligned}$$

where α is *maximal monotone*, or in the case of a measurable family $\{\alpha(x; \cdot) : x \in \Omega\}$, each is a maximal monotone relation.

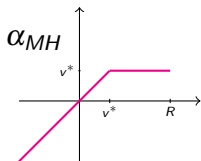
Examples: $\alpha(u) = \alpha_S, \alpha_{MH}, \underbrace{\alpha_E, \alpha_W}_{\text{Toy models}}, \alpha_{PM}$; Note: $\beta(v) = \alpha^{-1}$

Toy models



Stefan free-boundary problem:

$$\alpha_S(u) = u_- + (u - 1)_+$$

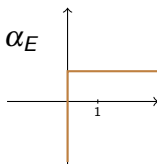


Methane hydrate (our problem):

$$\alpha_{MH}(x; \cdot) = (u - v^*(x))_- + v^*(x), \quad u \leq R$$

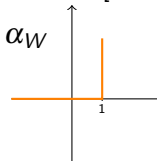
Porous medium equation:

$$\alpha = \alpha_{PM}(u) = |u|u^{m-1} \quad (m > 1 \text{ slow diffusion}, \quad 0 < m < 1 \text{ fast diffusion})$$



Elbow

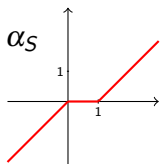
Showalter [1984]



Woble

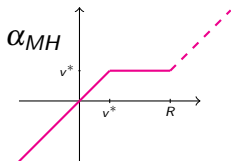
Examples: $\alpha(u) = \alpha_S, \alpha_{MH}, \underbrace{\alpha_E, \alpha_W}_{\text{Toy models}}, \alpha_{PM}$; Note: $\beta(v) = \alpha^{-1}$

Toy models



Stefan free-boundary problem:

$$\alpha_S(u) = u_- + (u - 1)_+$$

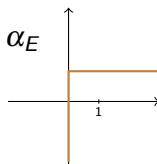


Methane hydrate (our problem):

$$\alpha_{MH}(x; \cdot) = (u - v^*(x))_- + v^*(x), \quad u \leq R$$

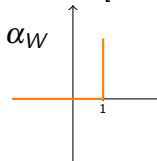
Porous medium equation:

$$\alpha = \alpha_{PM}(u) = |u|u^{m-1} \quad (m > 1 \text{ slow diffusion}, \quad 0 < m < 1 \text{ fast diffusion})$$



Elbow

Showalter [1984]



Woble

Summary of theoretical results for α_{MH}

- For the single graph case, we represent the non-linearity as a subgradient, and prove a useful comparison principle, which allows to extend the graph of $\beta = \alpha^{-1}$ to one which is affine bounded. Optimal regularity results follow.
 - ▶ Properties of u, v are the same as those for the Stefan problem.

Summary of theoretical results for α_{MH}

- For the single graph case, we represent the non-linearity as a subgradient, and prove a useful comparison principle, which allows to extend the graph of $\beta = \alpha^{-1}$ to one which is affine bounded. Optimal regularity results follow.
 - ▶ Properties of u, v are the same as those for the Stefan problem.
- We extend existing theory for *porous medium equation* to cover the case of a measurable family of graphs in order to show well-posedness.
 - ▶ Based on a normal convex integrand construction.

Details in Gibson, Medina, Peszyska, and Showalter [2013]

FE formulation for $\alpha = \alpha_{MH}(x, \cdot)$

First, we apply fully implicit time stepping.

Let $\mathcal{V}_h \subset \mathcal{V}$ be the finite element space of continuous piecewise linears on triangulation of Ω .

Find $v_h^n \in \mathcal{V}_h$ at t_n ($n > 0$)

$$\begin{cases} (u_h^n, \psi) + \tau(\nabla v_h^n, \nabla \psi) = (u_h^{n-1}, \psi) \\ u_h^n \in \beta(v_h^n) \\ (u_h^0, \psi) := (u_0, \psi), \forall \psi \in \mathcal{V}_h \end{cases}$$

FE formulation for $\alpha = \alpha_{MH}(x, \cdot)$

First, we apply fully implicit time stepping.

Let $\mathcal{V}_h \subset \mathcal{V}$ be the finite element space of continuous piecewise linears on triangulation of Ω .

Find $v_h^n \in \mathcal{V}_h$ at t_n ($n > 0$)

$$\begin{cases} (u_h^n, \psi) + \tau(\nabla v_h^n, \nabla \psi) = (u_h^{n-1}, \psi) \\ u_h^n \in \beta(v_h^n) \\ (u_h^0, \psi) := (u_0, \psi), \forall \psi \in \mathcal{V}_h \end{cases}$$

Let

M : mass matrix

K : stiffness matrix

$$v_h^n \approx \mathbf{v}^n \in \mathbb{R}^M$$

$$u_h^n \approx \mathbf{u}^n \in \mathbb{R}^M$$

$$\mathbf{M}\mathbf{u}^n + \tau\mathbf{K}\mathbf{v}^n = \mathbf{M}\mathbf{u}^{n-1}$$

FE formulation for $\alpha = \alpha_{MH}(x, \cdot)$

First, we apply fully implicit time stepping.

Let $\mathcal{V}_h \subset \mathcal{V}$ be the finite element space of continuous piecewise linears on triangulation of Ω .

Find $v_h^n \in \mathcal{V}_h$ at t_n ($n > 0$)

$$\begin{cases} (u_h^n, \psi) + \tau(\nabla v_h^n, \nabla \psi) = (u_h^{n-1}, \psi) \\ u_h^n \in \beta(v_h^n) \\ (u_h^0, \psi) := (u_0, \psi), \forall \psi \in \mathcal{V}_h \end{cases}$$

Let

M : mass matrix

K : stiffness matrix

$$v_h^n \approx \mathbf{v}^n \in \mathbb{R}^M$$

$$u_h^n \approx \mathbf{u}^n \in \mathbb{R}^M$$

$$\mathbf{M}\mathbf{u}^n + \tau\mathbf{K}\mathbf{v}^n = \mathbf{M}\mathbf{u}^{n-1}$$

Mass-lumping allows

$$\mathbf{A}_h = \mathbf{M}^{-1}\mathbf{K}$$

Fully discrete scheme

$$\begin{cases} \mathbf{u}^n + \tau\mathbf{A}_h\mathbf{v}^n = \mathbf{u}^{n-1} \\ \langle v_j^n, u_j^n \rangle \in \beta(x_j; \cdot) := \beta_j(\cdot) \end{cases}$$

where the constraint is applied point-wise.

Lemma [N.Gibson, P. Medina, M. Peszynska, R.E. Showalter]

For every $n > 0$ there is a unique solution $\mathbf{v}^n \in \mathbb{R}^M$ of the discrete problem for $\beta = \beta_{MH}(x; \cdot)$ it is the unique minimizer of the appropriate functional $\Psi(\mathbf{v})$ for which the discrete problem is the Euler-Lagrange condition.

Lemma [N.Gibson, P. Medina, M. Peszynska, R.E. Showalter]

For every $n > 0$ there is a unique solution $\mathbf{v}^n \in \mathbb{R}^M$ of the discrete problem for $\beta = \beta_{MH}(x; \cdot)$ it is the unique minimizer of the appropriate functional $\Psi(\mathbf{v})$ for which the discrete problem is the Euler-Lagrange condition.

Corollary

The discrete scheme is uniquely solvable for each of $\beta = \beta_{MH}(x, \cdot), \beta_S, \beta_E, \beta_W$.

Nonlinear algebraic subproblem

- Newton-type solvers have difficulties near singularities; may not be defined for multi-valued operators.

Nonlinear algebraic subproblem

- Newton-type solvers have difficulties near singularities; may not be defined for multi-valued operators.
- Relaxation solvers require a number of iterations proportional to the number of degrees of freedom.

Nonlinear algebraic subproblem

- Newton-type solvers have difficulties near singularities; may not be defined for multi-valued operators.
- Relaxation solvers require a number of iterations proportional to the number of degrees of freedom.
- In each, optimal convergence results are close to $O(h)$ for v and $O(h^{\frac{1}{2}})$ for u in $L^2(Q)$.

Nonlinear algebraic subproblem

- Newton-type solvers have difficulties near singularities; may not be defined for multi-valued operators.
- Relaxation solvers require a number of iterations proportional to the number of degrees of freedom.
- In each, optimal convergence results are close to $O(h)$ for v and $O(h^{\frac{1}{2}})$ for u in $L^2(Q)$.
- We propose a scheme which does not require regularization and can be applied when neither α nor β are functions.

Nonlinear algebraic subproblem

- Newton-type solvers have difficulties near singularities; may not be defined for multi-valued operators.
- Relaxation solvers require a number of iterations proportional to the number of degrees of freedom.
- In each, optimal convergence results are close to $O(h)$ for v and $O(h^{\frac{1}{2}})$ for u in $L^2(Q)$.
- We propose a scheme which does not require regularization and can be applied when neither α nor β are functions.
- The method also applies when constraints are parameterized by x .

Nonlinear Complementarity Problem (NCP)

We represent

$$\langle v, u \rangle \in \beta$$

as an **NCP-function**

$$\phi(u, v) = 0.$$

Nonlinear Complementarity Problem (NCP)

We represent

$$\langle v, u \rangle \in \beta$$

as an **NCP-function**

$$\phi(u, v) = 0.$$

For example,

$$\langle v, u \rangle \in \beta_{MH} \equiv \phi_{MH}(u, v) := \min(u - v, v^*(x) - v) = 0.$$

Nonlinear Complementarity Problem (NCP)

We represent

$$\langle v, u \rangle \in \beta$$

as an **NCP-function**

$$\phi(u, v) = 0.$$

For example,

$$\langle v, u \rangle \in \beta_{MH} \equiv \phi_{MH}(u, v) := \min(u - v, v^*(x) - v) = 0.$$

Similarly,

$$\langle v, u \rangle \in \beta_E \equiv \phi_E(u, v) := \min(u, 1 - v) = 0,$$

$$\langle v, u \rangle \in \beta_W \equiv \phi_W(u, v) := \min(1 - u, v) = 0,$$

$$\langle v, u \rangle \in \beta_S \equiv \phi_S(u, v) := u - v - \max(0, \min(u, 1)) = 0.$$

Semismooth Newton solver

Problem solved at every time step becomes

$$\begin{cases} \mathbf{u} + \tau \mathbf{A}_h \mathbf{v} = \mathbf{b} \\ \min(u_j - v_j, v^*(x_j) - v_j) = 0, \forall j \end{cases}$$

Semismooth Newton solver

Problem solved at every time step becomes

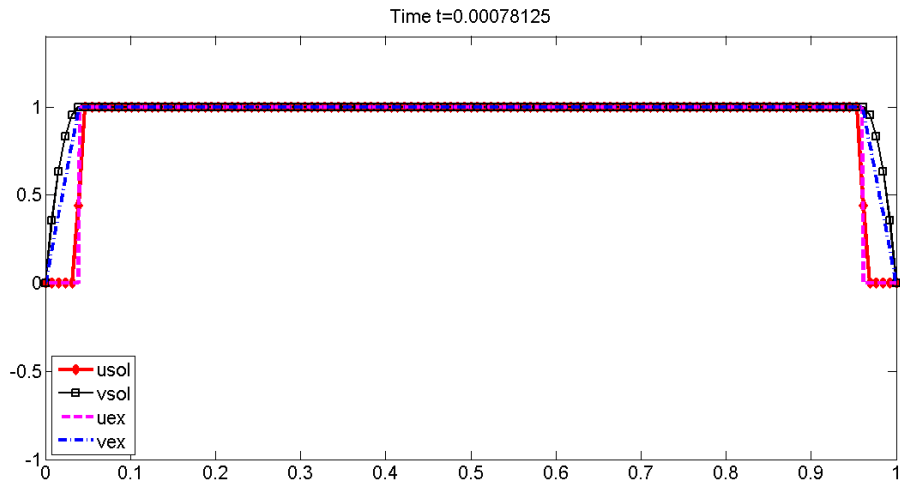
$$\begin{cases} \mathbf{u} + \tau \mathbf{A}_h \mathbf{v} = \mathbf{b} \\ \min(u_j - v_j, v^*(x_j) - v_j) = 0, \forall j \end{cases}$$

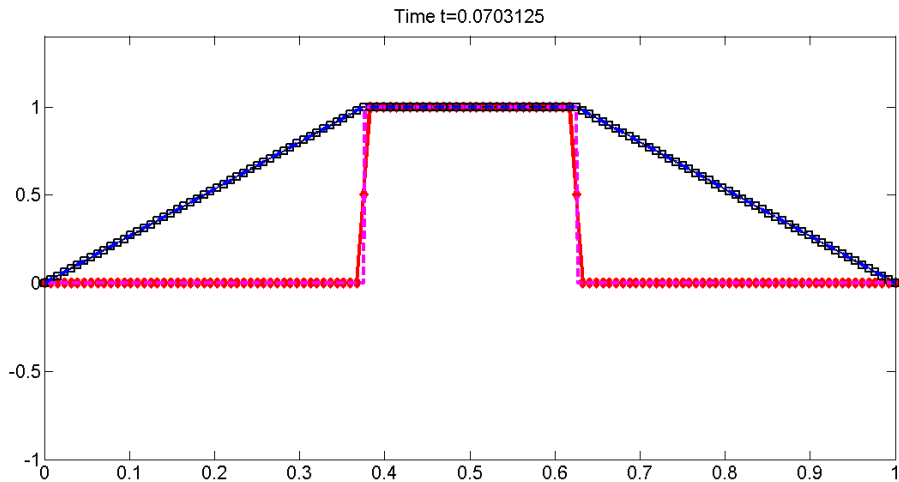
It can be shown that

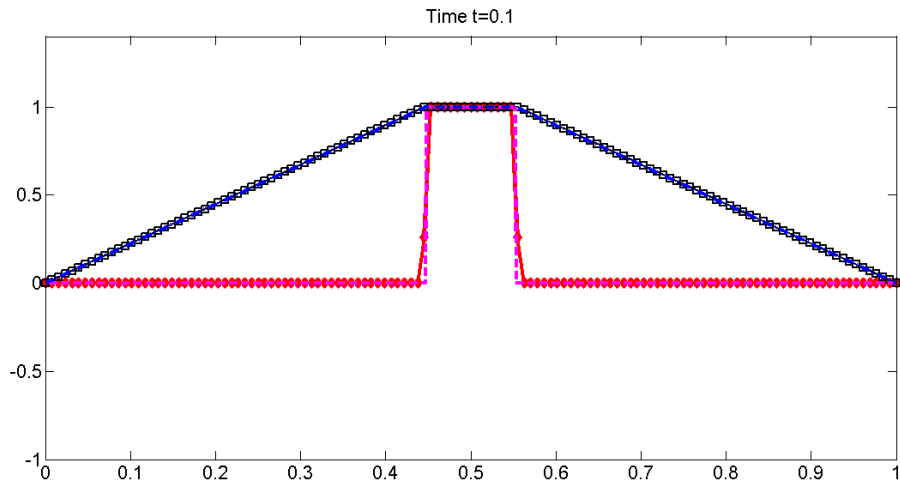
- each of $\phi_{MH}, \phi_E, \phi_W, \phi_S$ is semi-smooth
- the Jacobian is never singular

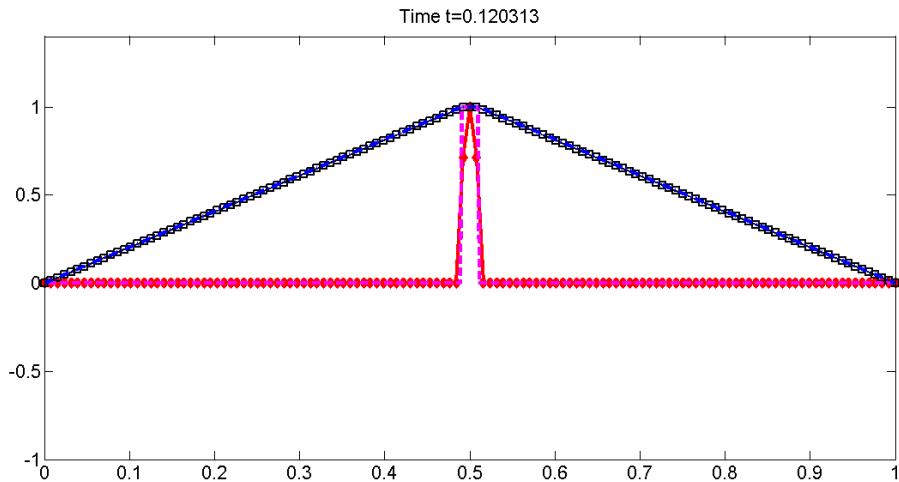
Semismooth Newton converges *superlinearly* for these NCC problems

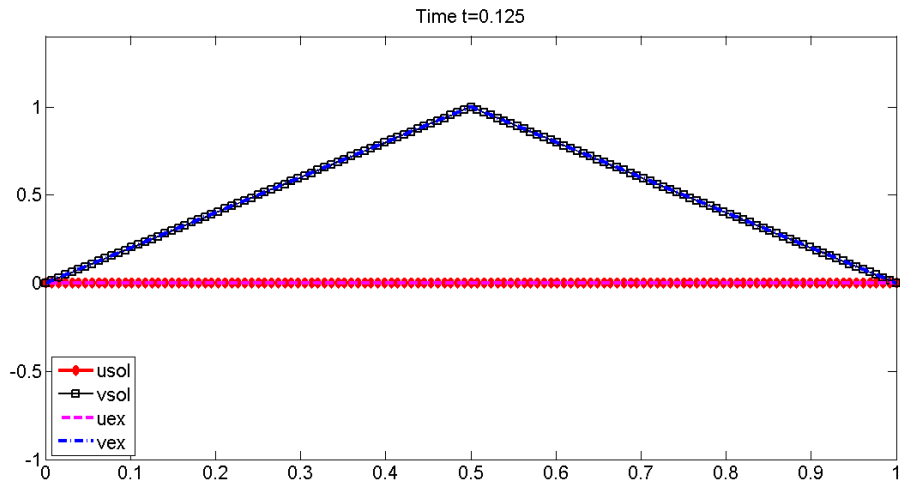
Ben Gharbia, Gilbert, and Jaffre [2011]; Ulbrich [2011]

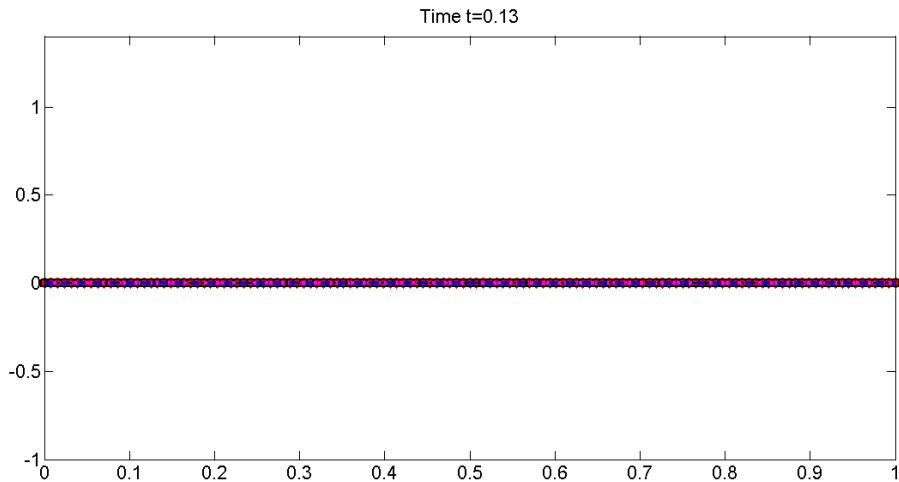
$\alpha = \alpha_E$. Toy model (Showalter [1984]) Frame I

$\alpha = \alpha_E$. Toy model (Showalter [1984]) Frame II

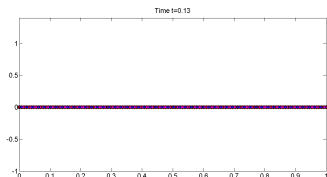
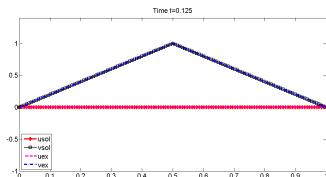
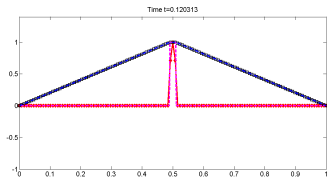
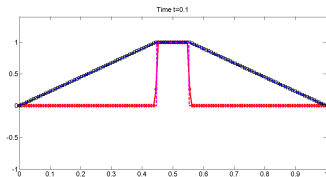
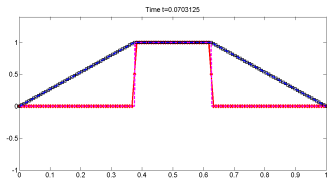
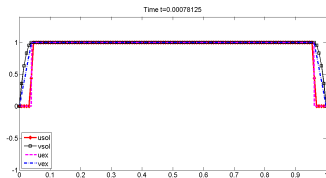
$\alpha = \alpha_E$. Toy model (Showalter [1984]) Frame III

$\alpha = \alpha_E$. Toy model (Showalter [1984]) Frame IV

$\alpha = \alpha_E$. Toy model (Showalter [1984]) Frame V

$\alpha = \alpha_E$. Toy model (Showalter [1984]) Frame VI

$\alpha = \alpha_E$. Toy model (Showalter [1984])



Convergence in u and v for α_E Using $\tau = \frac{h}{10}, \frac{h}{100}, h^2$

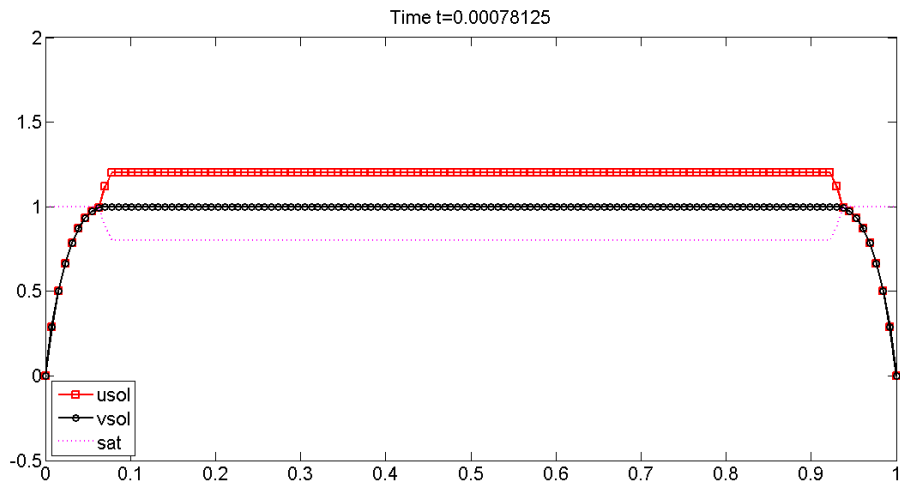
			error	rate	error	rate	error	rate
$1/h$	$1/\tau$	N_{it}	$e_{u,2}$	$r_{u,2}$	$e_{v,2}$	$r_{v,2}$	e_q	r_q
256	2560	2	1.03e-02	0.540	1.19e-03	0.785	6.40e-04	1.073
512	5120	2	6.81e-03	0.601	6.73e-04	0.828	3.00e-04	1.094
128	12800	2	1.47e-02	0.546	1.23e-03	0.966	1.48e-03	1.016
256	25600	2	9.69e-03	0.602	6.29e-04	0.961	7.19e-04	1.040
32	1024	2	2.90e-02	0.516	5.25e-03	0.945	5.42e-03	0.800
64	4096	2	1.93e-02	0.591	2.62e-03	1.003	2.78e-03	0.964

(with quasi-norm: $\sum_n \tau \int_{\Omega} |u - u_h^n| |v - v_h^n| dx$, Ebmeyer and Liu [2008]).

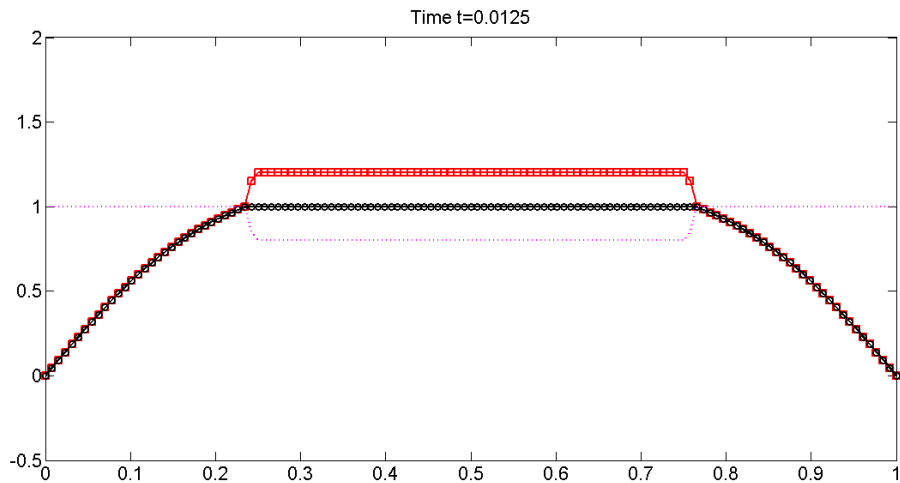
Observed rates

$$e_{u,2} \approx O(h^{1/2}), \quad e_{v,2} \approx O(h), \quad e_q \approx O(h)$$

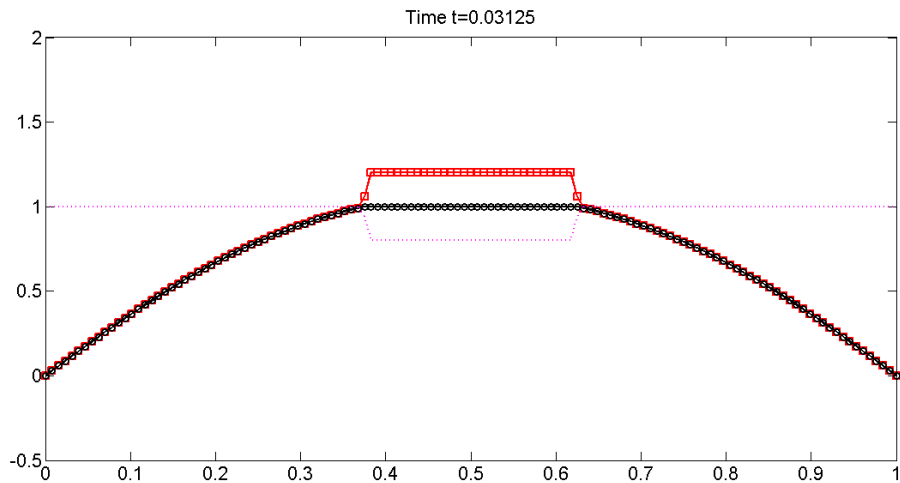
$\alpha = \alpha_{MH}$, $v_{\max}^* \equiv 1$. No analytical solution. Frame 1



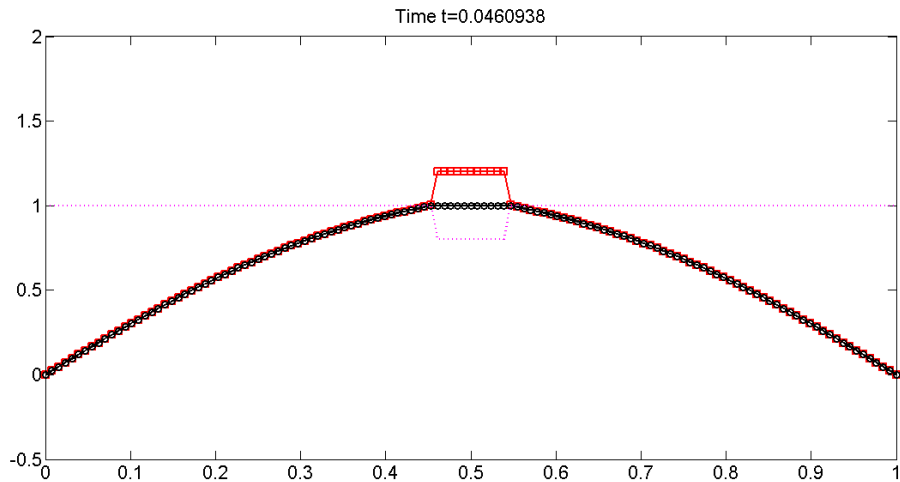
$\alpha = \alpha_{MH}$, $v_{\max}^* \equiv 1$. No analytical solution. Frame II



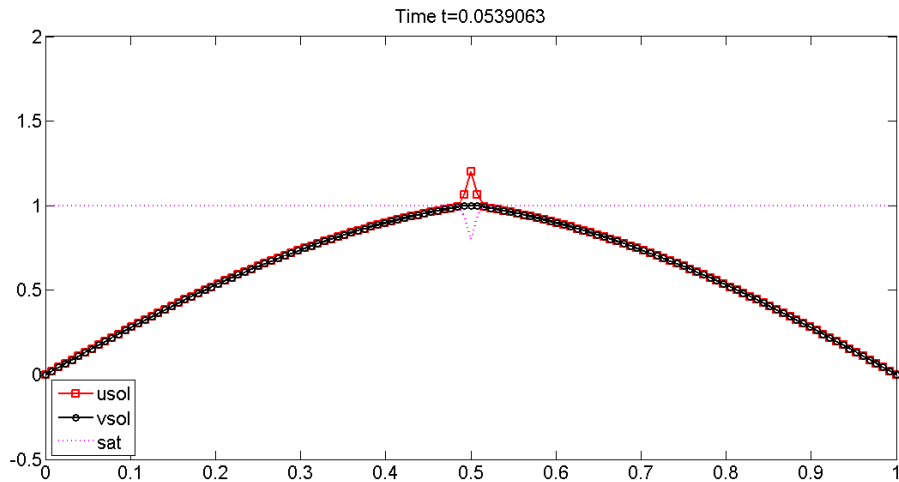
$\alpha = \alpha_{MH}$, $v_{\max}^* \equiv 1$. No analytical solution. Frame III



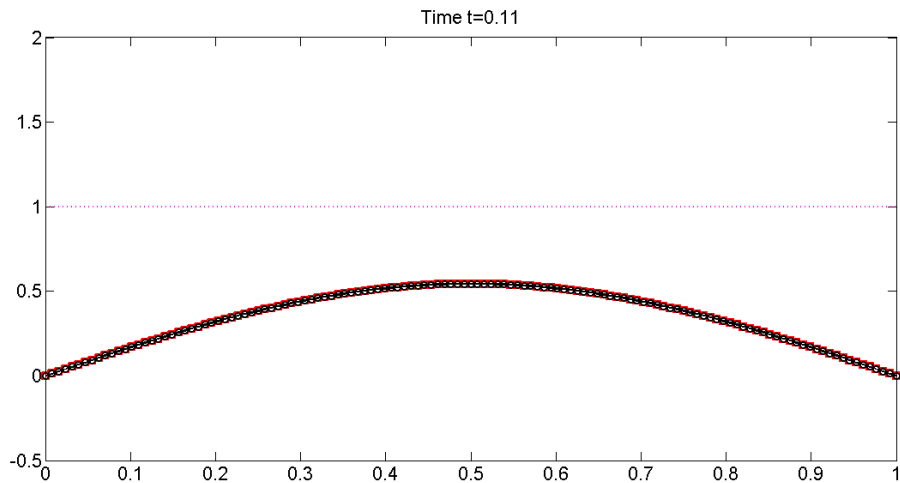
$\alpha = \alpha_{MH}$, $v_{\max}^* \equiv 1$. No analytical solution. Frame IV



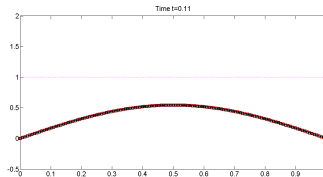
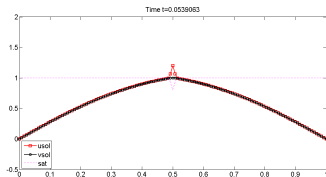
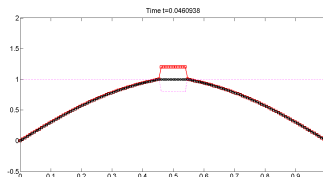
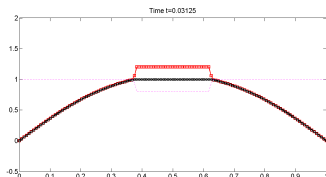
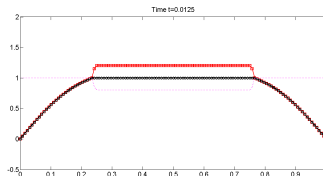
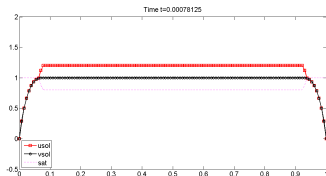
$\alpha = \alpha_{MH}$, $v_{\max}^* \equiv 1$. No analytical solution. Frame V



$\alpha = \alpha_{MH}, v_{\max}^* \equiv 1$. No analytical solution. Frame VI



$\alpha = \alpha_{MH}$, $v_{\max}^* \equiv 1$. No analytical solution.



Convergence in u and v for $\alpha_{MH}, v_{\max}^* \equiv 1$.

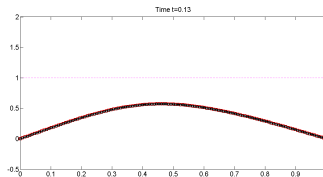
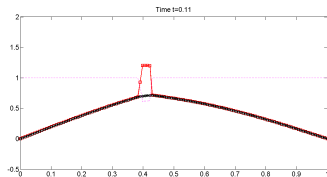
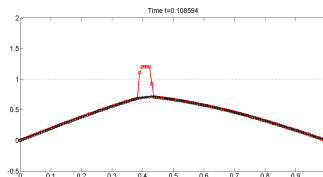
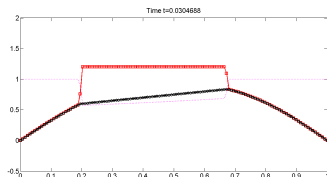
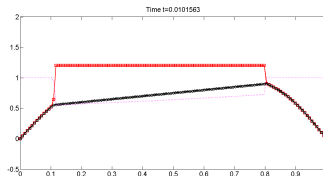
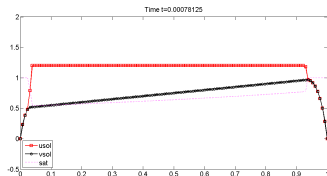
Using $\tau = \frac{h}{10}, \frac{h}{100}, h^2$

$1/h$	$1/\tau$	N_{it}	$e_{u,2}$	$r_{u,2}$	$e_{v,2}$	$r_{v,2}$	e_q	r_q
256	2560	2	2.14e-03	0.560	8.50e-04	0.760	8.62e-04	0.768
512	5120	2	1.39e-03	0.623	4.86e-04	0.806	4.92e-04	0.810
128	12800	2	1.98e-03	0.559	2.17e-04	0.682	3.33e-04	0.855
256	25600	2	1.31e-03	0.603	1.31e-04	0.725	1.81e-04	0.883
32	1024	2	4.39e-03	0.833	1.31e-03	1.396	1.55e-03	1.287
64	4096	2	2.69e-03	0.705	4.88e-04	1.421	6.57e-04	1.239

Observed rates

$$e_{u,2} \approx O(h^{1/2}), \quad e_{v,2} \approx O(h), \quad e_q \approx O(h)$$

$$\alpha = \alpha_{MH}(x; \cdot), v_{\max}^*(x) = (1+x)/2$$



Convergence rates in u and v for $\alpha = \alpha_{MH}(x; \cdot)$, $v_{\max}^*(x) = (1+x)/2$

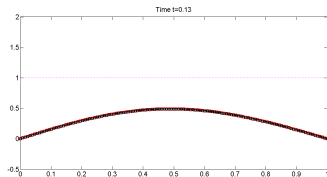
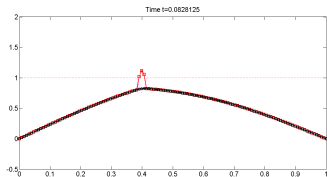
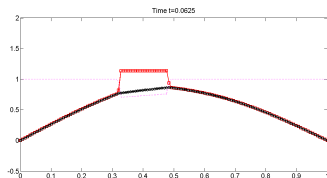
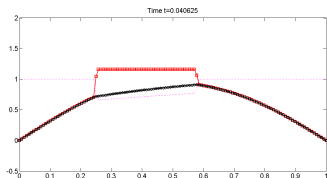
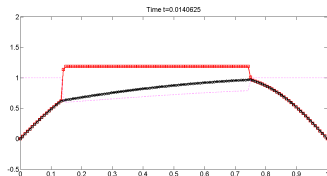
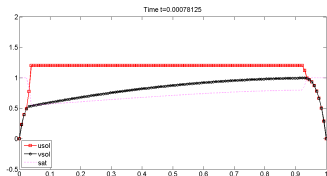
Using $\tau = \frac{h}{10}, \frac{h}{100}, h^2$

$1/h$	$1/\tau$	N_{it}	$e_{u,2}$	$r_{u,2}$	$e_{v,2}$	$r_{v,2}$	e_q	r_q
256	2560	2	5.22e-03	0.551	6.74e-04	0.762	7.27e-04	0.798
512	5120	2	3.44e-03	0.602	3.84e-04	0.810	4.07e-04	0.838
128	12800	2	7.11e-03	0.545	2.47e-04	1.039	6.94e-04	0.986
256	25600	2	4.69e-03	0.601	1.29e-04	0.941	3.45e-04	1.010
32	1024	2	1.43e-02	0.622	1.30e-03	1.262	2.57e-03	0.942
64	4096	2	9.35e-03	0.612	5.52e-04	1.236	1.29e-03	0.997

Observed rates

$$e_{u,2} \approx O(h^{1/2}), \quad e_{v,2} \approx O(h), \quad e_q \approx O(h)$$

$$\alpha = \alpha_{MH}(x; \cdot), v_{\max}^*(x) = (1 + 2x - x^2)/2$$



Convergence in u and v for $\alpha = \alpha_{MH}(x; \cdot)$, $v_{\max}^*(x) = (1 + 2x - x^2)/2$

Using $\tau = \frac{h}{10}, \frac{h}{100}, h^2$

$1/h$	$1/\tau$	N_{it}	$e_{u,2}$	$r_{u,2}$	$e_{v,2}$	$r_{v,2}$	e_q	r_q
256	2560	2	3.45e-03	0.561	6.77e-04	0.763	7.07e-04	0.785
512	5120	2	2.27e-03	0.605	3.86e-04	0.811	3.98e-04	0.827
128	12800	2	4.50e-03	0.554	2.19e-04	0.990	5.33e-04	0.967
256	25600	2	2.96e-03	0.604	1.19e-04	0.875	2.68e-04	0.995
32	1024	2	9.18e-03	0.636	1.18e-03	1.330	1.98e-03	1.013
64	4096	2	5.98e-03	0.619	4.86e-04	1.280	9.88e-04	1.000

Observed rates

$$e_{u,2} \approx O(h^{1/2}), \quad e_{v,2} \approx O(h), \quad e_q \approx O(h)$$

Convergence in S for $\alpha = \alpha_{MH}$

Using $\tau = \frac{h}{100}$

$1/h$	constant		affine		non-affine	
	$e_{S,2}$	$r_{S,2}$	$e_{S,2}$	$r_{S,2}$	$e_{S,2}$	$r_{S,2}$
64	2.91e-03	0.537	7.89e-03	0.519	5.27e-03	0.525
128	1.97e-03	0.559	5.41e-03	0.546	3.58e-03	0.556
256	1.30e-03	0.602	3.56e-03	0.600	2.36e-03	0.603

Observed rates

$$e_{S,2} \approx O(h^{1/2})$$

(Similar to rates in u .)

Summary

- We described a solubility constrained methane hydrate model.

Summary

- We described a solubility constrained methane hydrate model.
- We reformulated it into an abstract evolution equation constrained by parameter-dependent families of graphs.

Summary

- We described a solubility constrained methane hydrate model.
- We reformulated it into an abstract evolution equation constrained by parameter-dependent families of graphs.
- One can extend monotone operator theory to case of measurable family of graphs to show well-posedness.

Summary

- We described a solubility constrained methane hydrate model.
- We reformulated it into an abstract evolution equation constrained by parameter-dependent families of graphs.
- One can extend monotone operator theory to case of measurable family of graphs to show well-posedness.
- Regularity of solutions is the same as in the Stefan problem, at least in the single graph case.

Summary

- We described a solubility constrained methane hydrate model.
- We reformulated it into an abstract evolution equation constrained by parameter-dependent families of graphs.
- One can extend monotone operator theory to case of measurable family of graphs to show well-posedness.
- Regularity of solutions is the same as in the Stefan problem, at least in the single graph case.
- We have proposed a numerical scheme which applies semismooth Newton to complementarity conditions.

Summary

- We described a solubility constrained methane hydrate model.
- We reformulated it into an abstract evolution equation constrained by parameter-dependent families of graphs.
- One can extend monotone operator theory to case of measurable family of graphs to show well-posedness.
- Regularity of solutions is the same as in the Stefan problem, at least in the single graph case.
- We have proposed a numerical scheme which applies semismooth Newton to complementarity conditions.
- Semismooth Newton solver requires mesh independent iterations.

Summary

- We described a solubility constrained methane hydrate model.
- We reformulated it into an abstract evolution equation constrained by parameter-dependent families of graphs.
- One can extend monotone operator theory to case of measurable family of graphs to show well-posedness.
- Regularity of solutions is the same as in the Stefan problem, at least in the single graph case.
- We have proposed a numerical scheme which applies semismooth Newton to complementarity conditions.
- Semismooth Newton solver requires mesh independent iterations.
- Convergence rates for examples agree with optimal results for the Stefan problem.

Summary

- We described a solubility constrained methane hydrate model.
- We reformulated it into an abstract evolution equation constrained by parameter-dependent families of graphs.
- One can extend monotone operator theory to case of measurable family of graphs to show well-posedness.
- Regularity of solutions is the same as in the Stefan problem, at least in the single graph case.
- We have proposed a numerical scheme which applies semismooth Newton to complementarity conditions.
- Semismooth Newton solver requires mesh independent iterations.
- Convergence rates for examples agree with optimal results for the Stefan problem.
- Incidentally discovered a semismooth Newton method for the Stefan problem.

Future work

- Implementation and convergence studies for the gas zone.

Future work

- Implementation and convergence studies for the gas zone.
- Include salinity as unknown.

Future work

- Implementation and convergence studies for the gas zone.
- Include salinity as unknown.
- Semi-implicit time stepping.

References

- I. Ben Gharbia, J. C. Gilbert, and J. Jaffre. Nonlinear complementarity constraints for two-phase flow in porous media with gas phase appearance and disappearance. *Numerical approximation of hysteresis problems*, 2011.
- C. Ebmeyer and W. B. Liu. Finite element approximation of the fast diffusion and the porous medium equations. *SIAM J. Numer. Anal.*, 46(5):2393–2410, 2008.
- N. L. Gibson, F. P. Medina, M. Peszynska, and R. E. Showalter. Evolution of phase transitions in methane hydrate. *To appear in Journal of Mathematical Analysis and Applications*, 2013.
- R. H. Nochetto and C. Verdi. Approximation of degenerate parabolic problems using numerical integration. *SIAM J. Numer. Anal.*, 25(4):784–814, 1988.
- J. Rulla. Error analysis for implicit approximations to solutions to Cauchy problems. *SIAM J. Numer. Anal.*, 33(1):68–87, 1996.
- R. E. Showalter. A singular quasilinear diffusion equation in L^1 . *J. Math. Soc. Japan*, 36(2):177–189, 1984.
- R. E. Showalter. *Monotone operators in Banach space and nonlinear partial differential equations*, volume 49 of *Mathematical Surveys and Monographs*. American Mathematical Society, Providence, RI, 1997.
- Michael Ulbrich. *Semismooth Newton methods for variational inequalities and constrained optimization problems in function spaces*, volume 11 of *MOS-SIAM Series on Optimization*. SIAM, Philadelphia, PA, 2011.